

# Quantum Dynamics as Classical Dynamics and Quantum/Symplectic Computers

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# Outlook

- **Quantum-symplectic duality**
  - Schrodinger equation and symplectic mechanics
  - Unitary and symplectic evolution
  - Qubits and Symbits
  - Euler-Bernoulli Eq. Elasticity theory. Vibrating beam, plates.
- **Quantum/Symplectic computer**
  - Unitary group  $U(N)$  and symplectic group  $Sp(2N, \mathbb{R})$
  - NOT and  $\gamma(\text{NOT})$  gates
- **Conclusion**

# Quantum-symplectic duality

# Schrodinger equation and symplectic mechanics

- Main message: Schrodinger equation actually describes the classical symplectic Hamiltonian system.
- Euler-Bernoulli Eq. Elasticity theory. Vibrating beam, plates
- Quantum/Symplectic computer
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- Compare:  
wave-particle duality Einstein(photons)-de Broglie(matter waves),...

# Schrodinger equation and symplectic mechanics

- Hilbert space

$$\mathcal{H} = \mathbb{C}^{\mathbb{N}} \quad (1)$$

- Schrodinger equation (SE)

$$i\dot{\psi}_{\mathbf{a}} = \mathbf{H}_{\mathbf{ab}}\psi_{\mathbf{b}}$$

$$\mathbf{a}, \mathbf{b} = 1 \dots \mathbf{N}, \psi_{\mathbf{a}} \in \mathbb{C}, \quad \mathbf{H}_{\mathbf{ab}} = \overline{\mathbf{H}}_{\mathbf{ba}} \quad (2)$$

# Schrodinger equation in real variables

- $\psi_{\mathbf{a}} = \mathbf{q}_{\mathbf{a}} + i\mathbf{p}_{\mathbf{a}}$ , where  $q_a, p_a$  are real.
- $H_{ab} = K_{ab} + iL_{ab}$ ,
- $K_{ab} = K_{ba}$ ,  $L_{ab} = -L_{ba}$  matrices with real entries.
- Then SE:

$$\dot{\mathbf{q}}_{\mathbf{a}} = \mathbf{K}_{\mathbf{ab}}\mathbf{p}_{\mathbf{b}} + \mathbf{L}_{\mathbf{ab}}\mathbf{q}_{\mathbf{b}} \quad (3)$$

$$\dot{\mathbf{p}}_{\mathbf{a}} = -\mathbf{K}_{\mathbf{ab}}\mathbf{q}_{\mathbf{b}} + \mathbf{L}_{\mathbf{ab}}\mathbf{p}_{\mathbf{b}}(*) \quad (4)$$

# Schrodinger equation in real variables

- (\*) are classical symplectic Hamiltonian equations

$$\begin{aligned}\dot{\mathbf{q}}_{\mathbf{a}} &= \frac{\partial \mathbf{H}_{\text{sym}}}{\partial \mathbf{p}_{\mathbf{a}}}, \\ \dot{\mathbf{p}}_{\mathbf{a}} &= -\frac{\partial \mathbf{H}_{\text{sym}}}{\partial \mathbf{q}_{\mathbf{a}}},\end{aligned}\tag{5}$$

with the Hamiltonian

$$\mathbf{H}_{\text{sym}} = \frac{1}{2}(\mathbf{p}_{\mathbf{a}} \mathbf{K}_{\mathbf{ab}} \mathbf{p}_{\mathbf{b}} + \mathbf{q}_{\mathbf{a}} \mathbf{K}_{\mathbf{ab}} \mathbf{q}_{\mathbf{b}}) + \mathbf{p}_{\mathbf{a}} \mathbf{L}_{\mathbf{ab}} \mathbf{q}_{\mathbf{b}}$$

## Remark 1.

If the Hamiltonian equation in the Schrodinger equations depends on time

$$i\dot{\psi}_{\mathbf{a}} = \mathbf{H}_{\mathbf{ab}}(\mathbf{t})\psi_{\mathbf{b}} \quad (6)$$

with

$$\mathbf{H}_{\mathbf{ab}}(\mathbf{t}) = \mathbf{K}_{\mathbf{ab}}(\mathbf{t}) + i\mathbf{L}_{\mathbf{ab}}(\mathbf{t}), \quad (7)$$

then we get the following equation

$$\begin{aligned} \dot{\mathbf{q}}_{\mathbf{a}} &= \mathbf{K}_{\mathbf{ab}}(\mathbf{t})\mathbf{p}_{\mathbf{b}} + \mathbf{L}_{\mathbf{ab}}(\mathbf{t})\mathbf{q}_{\mathbf{b}} \\ \dot{\mathbf{p}}_{\mathbf{a}} &= -\mathbf{K}_{\mathbf{ab}}(\mathbf{t})\mathbf{q}_{\mathbf{b}} + \mathbf{L}_{\mathbf{ab}}(\mathbf{t})\mathbf{p}_{\mathbf{b}} \end{aligned}$$

the Hamiltonian equation

$$\mathbf{H}_{\text{sym}}(\mathbf{t}) = \frac{1}{2}(\mathbf{p}_{\mathbf{a}}\mathbf{K}_{\mathbf{ab}}(\mathbf{t})\mathbf{p}_{\mathbf{b}} + \mathbf{q}_{\mathbf{a}}\mathbf{K}_{\mathbf{ab}}(\mathbf{t})\mathbf{q}_{\mathbf{b}}) + \mathbf{p}_{\mathbf{a}}\mathbf{L}_{\mathbf{ab}}\mathbf{q}_{\mathbf{b}}.$$



## Schrodinger equation

$$i\dot{\psi} = \mathbf{H}\psi \quad (8)$$

where  $\mathbf{H} = \frac{d^2}{dx^2} + \mathbf{V}(\mathbf{x})$

$$\ddot{\Psi} + \mathbf{H}^2\Psi = 0 \quad (9)$$

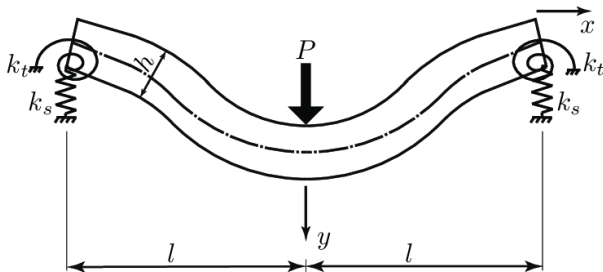
$$\Psi = \mathbf{q} + i\mathbf{p} \quad (10)$$

$$\ddot{\mathbf{q}} + \mathbf{H}^2\mathbf{q} = 0, \quad \mathbf{q} = \mathbf{q}(\mathbf{x}, \mathbf{t}) \quad (11)$$

# Elasticity theory

- Euler-Bernoulli vibrating beam eq.

$$\rho \ddot{W} + \frac{d^2}{dx^2} (\mathbf{K}(x) \frac{d^2}{dx^2} W) = 0 \quad (12)$$



# Vibrating plates

$$\rho \ddot{\mathbf{y}} + \mathbf{K}_1 \Delta \mathbf{y} + \mathbf{K}_2 \Delta \Delta \mathbf{y} = 0 \quad (13)$$

- Euler-Bernoulli 1650 year beam
- Eiffel Tower

$$\ddot{\mathbf{W}} + \frac{d^4}{dx^4} \mathbf{W} = 0 \quad (14)$$

- Schrodinger equation for free particle

$$i \dot{\Psi} = -\frac{d^2}{dx^2} \Psi \quad (15)$$

# Symplectic space

- Symplectic space is a pair  $(V, \omega)$
- $V$  - real vector space,  $\omega$  q-symmetric non-degenerate bilinear form
- Symplectic group is group of linear transformations  $(V)$  which preserves  $\omega$
- Example:  $\mathbb{R}(2N)$  there exist the standard symplectic structure which is given by the matrix

$$\begin{pmatrix} 0 & iN \\ -iN & 0 \end{pmatrix} \quad (16)$$

- The corresponding symplectic group is denoted  $\text{Sp}(2N, \mathbb{R})$

# Unitary and symplectic evolution

- Relation between representations of solutions of the Schrodinger/Hamiltonian equations in the complex unitary form and the real symplectic form.
- The solution of the N-component Schrodinger equation

$$\psi(\mathbf{t}) = \mathbf{U}_{\mathbf{t}} \psi(\mathbf{0}), \quad \psi(\mathbf{t}) = (\psi_1(\mathbf{t}), \dots, \psi_N(\mathbf{t})),$$

in term of real N-component vectors  $\psi(\mathbf{t}) = \mathbf{q}(\mathbf{t}) + i\mathbf{p}(\mathbf{t})$ .

# Unitary and symplectic evolution

- Representing the unitary matrix  $U_t$  in the form  $U_t = X_t + iY_t$ , where  $X_t$  and  $Y_t$  are real matrices, we have

$$\psi(t) = X_t q(0) - Y_t p(0) + i(Y_t q(0) + X_t p(0)),$$

i.e.  $q(t) = X_t q(0) - Y_t p(0), \quad p(t) = Y_t q(0) + X_t p(0) \quad (\#)$

- Relations (#) can be represented in the form

$$\begin{pmatrix} q(t) \\ p(t) \end{pmatrix} = S_t \begin{pmatrix} q(0) \\ p(0) \end{pmatrix}, \quad S_t = \begin{pmatrix} X_t & Y_t \\ -Y_t & X_t \end{pmatrix}$$

# Unitary and symplectic evolution

- The unitarity of matrix  $U_t$ ,

$$U_t^* U_t = (\mathbf{X}_t^T - i\mathbf{Y}_t^T)(\mathbf{X}_t + i\mathbf{Y}_t) \quad (17)$$

$$= \mathbf{X}_t^T \mathbf{X}_t + \mathbf{Y}_t^T \mathbf{Y}_t - i(\mathbf{Y}_t^T \mathbf{X}_t - \mathbf{X}_t^T \mathbf{Y}_t) = \mathbf{I} \quad (18)$$

means that

$$\mathbf{X}_t^T \mathbf{X}_t + \mathbf{Y}_t^T \mathbf{Y}_t = \mathbf{I},$$

$$\mathbf{Y}_t^T \mathbf{X}_t - \mathbf{X}_t^T \mathbf{Y}_t = \mathbf{0} \quad (*)$$

# Unitary and symplectic evolution

- Indeed,

$$\mathbf{S}_t^T \mathbf{J} \mathbf{S}_t = \begin{pmatrix} \mathbf{X}_t^T \mathbf{Y}_t - \mathbf{Y}_t^T \mathbf{X}_t & \mathbf{Y}_t^T \mathbf{Y}_t + \mathbf{X}_t^T \mathbf{X}_t \\ -\mathbf{Y}_t^T \mathbf{Y}_t - \mathbf{X}_t^T \mathbf{X}_t & -\mathbf{Y}_t^T \mathbf{X}_t + \mathbf{X}_t^T \mathbf{Y}_t \end{pmatrix}$$

and taking into account (\*) we get

$$\mathbf{S}_t^T \mathbf{J} \mathbf{S}_t = \begin{pmatrix} \mathbf{0} & \mathbf{I} \\ -\mathbf{I} & \mathbf{0} \end{pmatrix}, \quad \text{i.e. (**)}$$



# Qubits and Symbits

- In quantum computing, the fundamental unit of information is the qubit.
- The qubit is a two-level quantum system, mathematically represented by the complex Hilbert space  $\mathbb{C}^2$ .
- A general state of a qubit can be written as:

$$|\psi\rangle = \alpha|0\rangle + \beta|1\rangle, \quad \alpha, \beta \in \mathbb{C}, \quad |\alpha|^2 + |\beta|^2 = 1,$$

$|0\rangle$  and  $|1\rangle$  form an orthonormal basis for  $\mathbb{C}^2$

$$\alpha, \beta \in \mathbb{C}$$

# Qubits and Symbits

- In the symplectic computing framework, the analogous concept to a qubit is a symplectic bit (symbit).
- A symbit is represented by a two-dimensional real vector space  $\mathbb{R}^2$ .
- A general state of a symbit can be written as:

$$|\phi\rangle = \alpha|0\rangle + \beta|1\rangle, \quad \alpha, \beta \in \mathbb{R}, \quad \alpha^2 + \beta^2 = 1, \quad (19)$$

where  $|0\rangle$  and  $|1\rangle$  form an orthonormal basis for  $\mathbb{R}^2$ .

- The coefficients  $\alpha$  and  $\beta$  here are real numbers such that the total probability is equal to one.
- We can interpret  $\mathbb{R}^2$  as the phase plane for a dynamical system with coordinates  $p$  and  $q$ .

# Quantum mechanics over real number and Kahler space

- If  $\langle \cdot, \cdot \rangle$  is inner product in the complex Hilbert space then

$$[\langle \cdot, \cdot \rangle] = (\langle \cdot, \cdot \rangle) + i\omega \quad (20)$$

- where  $(\langle \cdot, \cdot \rangle)$  is positive defined and  $\omega$  skew-symmetric
- Kahler space  $(V, J, \omega)$ .  $V$  is a real vector space,  $J$  is a complex structure,  $J^2 = -1$  and  $\omega(\cdot, J\cdot)$  is positive defined.

# Tensor products of $N$ sbits

- Tensor products of  $N$  sbits:

$$\mathbb{R}^2 \otimes \mathbb{R}^2 \otimes \cdots \otimes \mathbb{R}^2 = \mathbb{R}^{2^N}. \quad (21)$$

This is equivalent to  $\mathbb{C}^N$ .

- On the other hand, the tensor product of  $N$  qubits is  $\mathbb{C}^{2^N}$ .
- Therefore, by using tensor products of sbits, one can obtain more general spaces than those obtained from tensor products of qubits.

# Tensor products of $N$ sbits

- Operations on symbits are performed using symplectic transformations, which are elements of the symplectic group  $Sp(2, \mathbb{R})$ .
- For instance, the simplest symplectic transformation in  $\mathbb{R}^2$  can be represented by a  $2 \times 2$  matrix  $S$  that preserves the symplectic form:

$$S^T J S = J, \quad \text{where} \quad J = \begin{pmatrix} 0 & 1 \\ -1 & 0 \end{pmatrix}.$$

An example of a symplectic transformation is the rotation matrix:

$$R(\theta) = \begin{pmatrix} \cos \theta & \sin \theta \\ -\sin \theta & \cos \theta \end{pmatrix},$$

which preserves the symplectic form and hence is a member of  $Sp(2, \mathbb{R})$ .

# Tensor products of $N$ sbits

- In summary, while qubits are the basic units of quantum information in  $\mathbb{C}^2$  and are manipulated using unitary transformations, symbits are the basic units of symplectic information in  $\mathbb{R}^2$  and are manipulated using symplectic transformations. This duality provides a bridge between quantum and symplectic computing, offering a new perspective on computational processes.

# Quantum/Symplectic computer

- Quantum computer is a sequence of unitary transformations (gates) and projectors (measurements)
- Symplectic computer is a sequence of symplectic transformations and projectors

$$U(N) = Sp(2N, \mathbb{R}) \cap O(2N, \mathbb{R}) \quad (22)$$

# Unitary and symplectic gates

- Unitary group  $U(N)$  and symplectic group  $Sp(2N, \mathbb{R})$

Let us show that unitary group  $U(N)$  is a subgroup of symplectic group  $Sp(2N, \mathbb{R})$ . Let us remind that a symplectic matrix  $S$  of symplectic group  $Sp(2N, \mathcal{R})$  satisfies the relation

$$S^T J S = J \quad (23)$$

where

$$J = \begin{pmatrix} 0 & I \\ -I & 0 \end{pmatrix}. \quad (24)$$



# Unitary and symplectic gates

- Unitary group  $U(N)$  and symplectic group  $Sp(2N, \mathbb{R})$

Let  $S$  is a  $2N \times 2N$  matrix in the form

$$S = \begin{pmatrix} A & B \\ C & D \end{pmatrix}$$

where  $A, B, C, D$  are  $N \times N$  matrices. The conditions on  $A, B, C$  and  $D$  are

- $A^T C, B^T D$  symmetric, and  $A^T D - C^T B = I$
- $AB^T, CD^T$  symmetric, and  $AD^T - BC^T = I$

# Unitary and symplectic gates

- Let us show that there is a canonical mapping  $\gamma$  of the unitary group  $U(N)$  to  $Sp(2N, \mathbb{R})$ . Let  $V \in U(N)$ , we present  $V$  in the form

$$V = X + iY, \quad \text{where } X, Y \text{ are } N \times N \text{ with real entries} \quad (25)$$

We define  $\gamma$  by the following formula

$$\gamma(V) = \gamma(X + iY) = \begin{pmatrix} X & Y \\ -Y & X \end{pmatrix} \quad (26)$$

One can check that the unitarity conditions  $VV^* = V^*V = 1$  lead to the conditions for the symplectic matrices.

## Remarks

- *Remark 2* One can also check that there is a **relation**  $\gamma(V_1)\gamma(V_2) = \gamma(V_1V_2)$ .
- *Remark 3.* **The Schrodinger equation in the form**

$$i\dot{\psi} = H\psi, \quad (27)$$

where

$$H = -\Delta + V(x) \quad (28)$$

$\Delta$  is the Laplace operator in  $\mathbb{R}$ , in the real formulation looks as follows

$$\dot{q} = Hp, \quad \dot{p} = -Hq. \quad (29)$$

Spectrum and scattering theory for in this approach are considered in [?]

# NOT and $\gamma$ (NOT) gates

- Quantum computation is a sequence of unitary matrices (gates) of the simple form. We put into correspondence into such unitary gate a corresponding symplectic gate using the  $\gamma$  defined by (26).
- See the simplest gate is NOT. It is a unitary matrix  $NOT = \begin{pmatrix} 0 & 1 \\ -1 & 0 \end{pmatrix}$  defined on the qubit  $\mathbb{C}^2$ . By using the mapping  $\gamma$  we define the corresponding symplectic gate

$$\gamma(NOT) = \begin{pmatrix} 0 & 1 & 0 & 0 \\ -1 & 0 & 0 & 0 \\ 0 & 0 & 0 & 1 \\ 0 & 0 & -1 & 0 \end{pmatrix} \quad (30)$$

One can construct the symplectic analog of the controlling CNOT gate.

# Conclusions

- Quantum-Symplectic duality
- Schrodinger equation is equivalent to Euler-Bernoulli equation for vibrating beam and elasticity theory.
- Symplectic computer is proposed